

AN INTEGRO-DIFFERENTIAL EQUATION WITHOUT CONTINUOUS SOLUTIONS

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ABSTRACT. We show an example of a non symmetric integro-differential equation of order α , for $\alpha \in (0, 1)$, for which Hölder estimates do not hold even though the kernels are comparable to the fractional Laplacian.

1. INTRODUCTION

We are concerned with an integro-differential equation of the usual form

$$(1) \quad \int_{\mathbb{R}^n} (u(x+y) - u(x) - y \cdot \nabla u(x) \chi_{B_1}(y)) K(x, y) dy = f(x) \text{ in } B_1.$$

The purpose of this article is to show an example of a kernel $K(x, y)$ satisfying, for some $\alpha \in (0, 1)$,

$$(2) \quad \frac{\lambda}{|y|^{n+\alpha}} \leq K(x, y) \leq \frac{\Lambda}{|y|^{n+\alpha}}.$$

and a bounded function f , for which the solution u of (1) does not satisfy any modulus of continuity a priori in terms of $\|u\|_{L^\infty}$ and $\|f\|_{L^\infty}$.

The key of our example is that we do not make the symmetry assumption $K(x, y) = K(x, -y)$. The correction term $-y \cdot \nabla u(x) \chi_{B_1}(y)$ inside the integrand effectively creates a drift term. Since we take $\alpha < 1$, the regularization effect of the symmetric part of the integral does not compensate the effect of this implicit drift. Any modulus of continuity can thus be invalidated following a mechanism similar to that in [11].

When the kernel K satisfies the symmetry assumption $K(x, y) = K(x, -y)$, then the solutions to (1) satisfy a regularity estimate in Hölder spaces

$$\|u\|_{C^\gamma(B_{1/2})} \leq C (\|u\|_{L^\infty(\mathbb{R}^n)} + \|f\|_{L^\infty(B_1)}),$$

for some $\gamma > 0$. This estimate was obtained first by R. Bass and D. Levin [2]. Hölder estimates of this type have been a topic of high interest in recent years, with several results in this direction for different types of integral equations including [12], [1], [9], [3], [4], [6], and also results for parabolic integral equations like [7], [10] and [8].

In some cases, Hölder estimates hold for non symmetric kernels $K(x, y)$. That is the case of the results in [4], [6] and [8]. In those cases, for $\alpha < 1$, the equation has to be taken without the gradient correction term. That is, the equation is

$$\int_{\mathbb{R}^n} (u(x+y) - u(x)) K(x, y) dy = f(x) \text{ in } B_1.$$

Note that for $\alpha \in (0, 1)$, the left hand side makes sense for every function $u \in C^1$. While equation (1) is of the form that appears traditionally in the probability literature, it is better not to include the gradient correction term in this case. The result in this note shows that in fact, this correction term ruins any continuity estimate. There have been attempts to obtain Hölder continuity estimates for equations of this form, see for example [5].

Now we state our main result.

Theorem 1. *For any $\alpha \in (0, 1)$ and $0 < \lambda < \Lambda$, and any modulus of continuity η , there is a kernel $K(x, y)$ satisfying (2), and a function $u : \mathbb{R}^n \rightarrow \mathbb{R}$ such that*

- $u(x) \in [-1, 1]$ for all $x \in \mathbb{R}^n$.
- u solves the equation (1) with kernel $K(x, y)$ and $f(x) \equiv 0$.
- The function u does not obey the modulus of continuity η at the origin.

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Remark. The solution u constructed in the next section is continuous on \mathbb{R}^n , smooth on the set $\{x_1 \neq 0\}$, and solves (1) classically on $B_1 \setminus \{x_1 = 0\}$. On the line $\{x_1 = 0\}$, $\nabla u(x)$ does not exist, but the equation is satisfied in the viscosity sense, as there are no C^2 functions touching u from above or below. More precisely, let M^\pm be the extremal operators

$$M^+u(x) = \sup \left\{ \int_{\mathbb{R}^n} (u(x+y) - u(x) - y \cdot \nabla u(x) \chi_{B_1}(y)) K(y) dy : \frac{\lambda}{|y|^{n+\alpha}} \leq K(y) \leq \frac{\Lambda}{|y|^{n+\alpha}} \right\},$$

$$M^-u(x) = \inf \left\{ \int_{\mathbb{R}^n} (u(x+y) - u(x) - y \cdot \nabla u(x) \chi_{B_1}(y)) K(y) dy : \frac{\lambda}{|y|^{n+\alpha}} \leq K(y) \leq \frac{\Lambda}{|y|^{n+\alpha}} \right\}.$$

Then, the function u from Theorem 1 satisfies $M^+u \geq 0$ and $M^-u \leq 0$ in the viscosity sense in B_1 .

See [5] for a more explicit expression of the operators M^+ and M^- .

2. THE PROOF

We give the value of the kernel $K(x, y)$ below. The function u will depend on the variable x_1 only. Thus, our example is essentially one-dimensional, and we will use the notations $u(x)$ and $u(x_1)$ interchangeably. The strategy of the proof is as follows: first, construct a family of bounded, continuous functions u_r that approximate the discontinuous function $u_0(x_1) = \text{sgn}(x_1)$ as $r \rightarrow 0$. If $r > 0$ is chosen small enough, u_r will fail to admit a given modulus of continuity. Next, we will add a continuous, increasing function $v(x_1)$ to $u_r(x_1)$, with v independent of r , to ensure that the quantity $L(u_r + v)(x)$, with

$$(3) \quad Lu(x) = \int_{\mathbb{R}^n} (u(x+y) - u(x) - y \cdot \nabla u(x) \chi_{B_1}(y)) K(x, y) dy,$$

satisfies $-C_0 \leq L(u_r + v)(x) \leq C_0$ for $x \in B_1$, with C_0 independent of r . (See Lemma 1.) Since v is increasing in x_1 , the function $u_r + v$ will still break the given modulus of continuity. Finally, we will add another continuous, increasing function $w(x_1)$ to our solution (see Lemma 2), such that the sum $u_r(x_1) + v(x_1) + w(x_1)$ satisfies $L(u_r + v + w)(x_1) = 0$ for $x_1 \in [-1, 1]$, and $u_r + v + w$ will also break the given modulus of continuity.

For $0 < \alpha < 1$ and $0 < \lambda < \Lambda$, let us define the kernel $K(x, y)$ as follows: let

$$K_1(x, y) = \frac{(\lambda + \Lambda)/2}{|y|^{n+\alpha}},$$

$$K_2(x, y) = \text{sgn}(y_1) \frac{\Lambda - \lambda}{2} \chi_{B_1}(y),$$

$$K_3(x, y) = \text{sgn}(y_1) \frac{\Lambda - \lambda}{2} \frac{1}{|y|^{n+\alpha}} \chi_{\mathbb{R}^n \setminus B_1}(y).$$

Note that $K_1(x, y)$ is even in y , and $K_2(x, y)$ and $K_3(x, y)$ are odd. Our kernel is defined by

$$(4) \quad K(x, y) := K_1(x, y) + a(x)K_2(x, y) - c(x)K_3(x, y),$$

where $a(x)$ is to be chosen later. The condition (2) implies that we need $|a(x)| \leq 1$ and $|c(x)| \leq 1$ for all x . Let $b(x)$ be the drift vector, which is given by

$$b(x) = \int_{B_1} y K_3(x, y) dy = \frac{\Lambda - \lambda}{2} \int_{B_1} y \text{sgn}(y_1) dy = \frac{\Lambda - \lambda}{2} \frac{\omega_n}{2} e_1,$$

where ω_n is a constant depending on the dimension n only.

We define

$$(5) \quad \begin{aligned} L_i u(x) &= \int_{\mathbb{R}^n} (u(x+y) - u(x)) K_i(x, y) dy, \quad i = 1, 2, 3, \\ L_4 u(x) &= b(x) \cdot \nabla u(x). \end{aligned}$$

With this notation,

$$(6) \quad Lu(x) = L_1 u(x) + a(x)(L_2 u(x) - L_4 u(x)) - c(x)L_3 u(x).$$

Note that L_1 is a multiple of the usual fractional Laplacian: $L_1 u = -c_{n,\alpha}(-\Delta)^{\alpha/2} u$ for some constant $c_{n,\alpha} > 0$, and for u depending only on x_1 , $L_4 = C_1 \partial_{x_1}$, where $C_1 = \omega_n(\Lambda - \lambda)/4$. It is well known that the gradient and the fractional Laplacian have the following scaling: if $u_r(x) = u(x/r)$, then

$$(7) \quad L_1 u_r(x) = r^{-\alpha} [L_1 u](x/r) \text{ and } L_4 u_r(x) = r^{-1} [L_4 u](x/r),$$

and we will repeatedly make use of this. Furthermore, simple integral estimates show that $L_2u(x)$ and $L_3u(x)$ are bounded in B_1 for any function u that is bounded in \mathbb{R}^n .

The following lemma establishes a solution u to equation (1) with bounded right-hand side f , such that u breaks a given modulus of continuity.

Lemma 1. *For any $\alpha \in (0, 1)$, $0 < \lambda < \Lambda$, and modulus of continuity η at the origin, there exist a bounded function $u : \mathbb{R}^n \rightarrow \mathbb{R}$ and a kernel of the form (4), with $c(x) = 0$, such that*

$$(8) \quad -C_0 \leq Lu(x) \leq C_0, \quad x \in B_1,$$

with L as in (3), and u breaks η . The function u depends only on x_1 and is monotonically increasing. The constant C_0 and $\sup_{\mathbb{R}^n} |u|$ depend on α , λ , Λ , and n , but are independent of η .

Proof. Let $u_1(x_1)$ be any smooth, nondecreasing function such that $u_1(x_1) = -1$ for $x_1 \leq -1$ and $u_1(x_1) = 1$ for $x_1 \geq 1$, and define $u_r(x_1) = u_1(x_1/r)$. Choose $r > 0$ small enough that $\eta(r) < 1$. Then the oscillation $\omega_{u_r}(B_r) = \sup_{B_r} u_r(x) - \inf_{B_r} u_r(x) = 1 \geq \eta(r)$, so $u_r(x)$ breaks η .

The function u satisfying (8) will be of the form $u(x) = u_r(x_1) + v(x_1)$, with v another function depending only on x_1 and increasing in x_1 . To define v , we pick a small $\varepsilon > 0$ (such that $\alpha + \varepsilon < 1$) and set

$$v(x_1) = \begin{cases} -2^{1-\alpha-\varepsilon}, & x_1 < -2, \\ \operatorname{sgn}(x_1)|x_1|^{1-\alpha-\varepsilon}, & |x_1| \leq 2, \\ 2^{1-\alpha-\varepsilon}, & x_1 > 2. \end{cases}$$

Since v is nondecreasing in x_1 , the function $u(x) = u_r(x) + v(x)$ also breaks the modulus of continuity η .

We claim there exists $\delta > 0$, independent of r , such that $L_4u(x_1) \geq |L_1u(x_1)|$ for $|x_1| < \delta$. To establish this, we first estimate L_1u_1 . By symmetry, we have

$$(9) \quad L_1u_1(x_1) = c_0 \int_{\mathbb{R}} \frac{u_1(x_1 + y_1) - u_1(x_1)}{|y_1|^{1+\alpha}} dy_1,$$

where c_0 depends on α, λ, Λ , and n . For $x_1 > 1$, we have

$$(10) \quad |L_1u_1(x_1)| \leq c_0 \int_{-\infty}^{1-x_1} \frac{2}{|y_1|^{1+\alpha}} dy_1 = \frac{2c_0}{\alpha} |x_1 - 1|^{-\alpha} \leq C|x_1|^{-\alpha},$$

for some constant C , and similarly, $|L_1u_1(x_1)| \leq C|x_1|^{-\alpha}$ for $x_1 < -1$. For $|x_1| \leq 1$, we have

$$|L_1u_1(x_1)| \leq \int_{-\infty}^{-1-x_1} \frac{2}{|y_1|^{1+\alpha}} dy_1 + \int_{-1-x_1}^{1-x_1} \frac{u_1(x_1 + y_1) - u_1(x_1)}{|y_1|^{1+\alpha}} dy_1 + \int_{1-x_1}^{\infty} \frac{2}{|y_1|^{1+\alpha}} dy_1.$$

Similarly to (10), the first and third terms are bounded by $C|x_1|^{-\alpha}$, and since u_1 is smooth, the middle term is bounded by some constant, uniformly in $|x_1| \leq 1$. We conclude that $|L_1u_1(x_1)| \leq C|x_1|^{-\alpha}$ holds uniformly in $x_1 \in \mathbb{R}$, for some C . Combined with the scaling (7), this implies that for u_r ,

$$|L_1u_r(x_1)| = r^{-\alpha} |L_1u_1(x/r)| \leq C|x_1|^{-\alpha},$$

for some constant C independent of r .

Next, we estimate L_1v . Letting $\tilde{v}(x_1) = \operatorname{sgn}(x_1)|x_1|^{1-\alpha-\varepsilon}$, the homogeneity $\tilde{v}(\lambda x_1) = \lambda^{1-\alpha-\varepsilon}\tilde{v}(x_1)$ and the scaling (7) imply that $|L_1\tilde{v}(x_1)| = C|x_1|^{1-2\alpha-\varepsilon}$ for some C . For $|x_1| \leq 1$, we have

$$\begin{aligned} |L_1v(x_1)| &\leq |L_1\tilde{v}(x_1)| + |L_1(\tilde{v} - v)(x_1)| \\ &\leq C|x_1|^{1-2\alpha-\varepsilon} + \int_{\mathbb{R} \setminus [-2-x_1, 2-x_1]} \frac{|x_1 + y_1|^{1-\alpha-\varepsilon} - 2^{1-\alpha-\varepsilon}}{|y_1|^{1+\alpha}} dy_1 \\ &\leq C|x_1|^{1-2\alpha-\varepsilon} + \int_{-\infty}^{-2-x_1} \frac{C|y_1|^{1-\alpha-\varepsilon}}{|y_1|^{1+\alpha}} dy_1 + \int_{2-x_1}^{\infty} \frac{C|y_1|^{1-\alpha-\varepsilon}}{|y_1|^{1+\alpha}} dy_1 \\ &\leq C|x_1|^{1-2\alpha-\varepsilon} + C(|x_1 - 2|^{1-2\alpha-\varepsilon} + |x_1 + 2|^{1-2\alpha-\varepsilon}) \\ &\leq C|x_1|^{1-2\alpha-\varepsilon}, \end{aligned}$$

where C denotes a changing constant.

We have $L_4u_r(x_1) = \frac{C_1}{r}\chi_{[-r,r]}(x_1)$ and $L_4v(x_1) = (1-\alpha-\varepsilon)|x_1|^{-\alpha-\varepsilon}\chi_{[-2,2]}(x_1)$. Since $L_4u = L_4u_r + L_4v$ grows faster than $|L_1u| = |L_1u_r + L_1v|$ as $x_1 \rightarrow 0$, there is a $\delta > 0$ such that $L_4u(x_1) \geq |L_1u(x_1)|$ when $|x_1| < \delta$, as required. This δ does not depend on r .

We now choose

$$(11) \quad a(x) = \begin{cases} \frac{L_1u(x)}{L_4u(x)}, & |x_1| < \delta, \\ 0, & |x_1| \geq \delta, \end{cases}$$

so that $Lu(x) = a(x)L_2u(x)$ on $|x_1| < \delta$. On $\delta \leq |x_1| \leq 1$, we have $Lu(x) = L_1u(x)$. By the above estimates, $L_1(u_r + v)$ is bounded for $\delta \leq |x_1| \leq 1$. Since L_2u is bounded for any bounded u , we conclude that $u(x) = u_r(x_1) + v(x_1)$ satisfies (8). \square

Note that the result of the previous lemma remains true for any choice of $c(x) \in [-1, 1]$, since $L_3u(x)$ is a bounded function, for any bounded u . A nonzero choice of $c(x)$ will be used to make the right hand side of the equation zero and obtain our main result.

In the next lemma, we find a function w such that $|L_1w|$ is small, L_3w is large, and L_2w and L_4w cancel each other. If we add w to the function $u_r + v$ from Lemma 1, these properties will allow us to choose $a(x)$ and $c(x)$ in (3) such that $L(u_r + v + w) = 0$.

Lemma 2. *For any constant $C_0 > 0$, there exists a bounded function w depending only on x_1 , and monotonically increasing in x_1 , satisfying*

$$(12) \quad L_3w(x) - |L_1w(x)| \geq C_0,$$

$$(13) \quad L_2w(x) - L_4w(x) = 0,$$

for all $x \in B_1$, where L_i are defined in (5).

Proof. Let $w_1(x_1)$ be defined by

$$w_1(x_1) = \begin{cases} -1, & x_1 < -1, \\ x_1, & |x_1| \leq 1, \\ 1, & x_1 > 1, \end{cases}$$

and let $w_K(x_1) = Kw_1(x_1/K)$, where $K > 2$ is a large number to be determined later. Since w_K is linear when $|x_1| \leq 1$, the identity $w_K(x+y) - w_K(x) - y \cdot \nabla w_K(x) = 0$ holds there, so we have $L_2w_K(x) - L_4w_K(x) = 0$ in B_1 .

Next, we claim that for K large enough, $|L_1w_K(x_1)|$ will be uniformly bounded by an arbitrarily small constant c for $|x_1| \leq 1$. By a direct computation, if $|x_1| \leq 1$, we have

$$\begin{aligned} L_1w_1(x_1) &= c_0 \left(\int_{-\infty}^{-1-x_1} \frac{-1-x_1}{|y_1|^{1+\alpha}} dy_1 + \int_{-1-x_1}^{1-x_1} \frac{y_1}{|y_1|^{1+\alpha}} dy_1 + \int_{1-x_1}^{\infty} \frac{1-x_1}{|y_1|^{1+\alpha}} dy_1 \right) \\ &= \frac{c_0}{\alpha(1-\alpha)} (|x_1 - 1|^{1-\alpha} - |x_1 + 1|^{1-\alpha}), \end{aligned}$$

with c_0 as in (9). Since $g(x_1) = L_1w_1(x_1)$ is differentiable at $x_1 = 0$, we have $|g(x_1)|/|x_1| \rightarrow C$ as $x_1 \rightarrow 0$, for some constant C . By the scaling (7), this implies

$$(14) \quad |L_1w_K(x_1)| = K^{1-\alpha} |L_1w_1(x_1/K)| \leq K^{1-\alpha} C |x_1/K| = CK^{-\alpha},$$

for $|x_1| \leq 1$, with C independent of K . Therefore, for K large enough, $|L_1w_K(x_1)| \leq c < 1$ for all $|x_1| \leq 1$ and a small constant c .

For $L_3 w_K$, since the integrand $(w_K(x_1 + y_1) - w_K(x_1)) \operatorname{sgn}(y_1) / |y|^{n+\alpha}$ is positive everywhere, we can write

$$\begin{aligned} L_3 w_K(x_1) &= \int_{\mathbb{R}^n \setminus B_1} \frac{w_K(x_1 + y_1) - w_K(x_1)}{|y|^{n+\alpha}} \operatorname{sgn}(y_1) dy_1 \\ &\geq \int_{B_{K-1} \setminus B_1} \frac{|y_1|}{|y|^{n+\alpha}} dy_1 \\ &= \int_1^{K-1} \int_{\mathbb{S}^{n-1}} \frac{\rho |\theta_1|}{\rho^{n+\alpha}} \rho^{n-1} d\theta d\rho \\ &= \frac{(K-1)^{1-\alpha} - 1}{1-\alpha} \int_{\mathbb{S}^{n-1}} |\theta_1| d\theta \\ &= C((K-1)^{1-\alpha} - 1). \end{aligned}$$

This lower bound holds uniformly for $|x_1| \leq 1$, so for K large enough, $\inf_{|x_1| \leq 1} w_K(x_1) > 1$. This implies we can choose K , depending on α, n, λ , and Λ , such that

$$L_3 w_K(x_1) - |L_1 w_K(x_1)| \geq 1 - c > 0, \quad |x_1| \leq 1.$$

Therefore, given $C_0 > 0$, we can choose a constant C such that

$$w(x_1) := C w_K(x_1)$$

satisfies (12) and (13). \square

We are now in a position to prove our result.

Proof of Theorem 1. Fix a modulus of continuity η . Let $\bar{u} = \bar{u}(x_1)$ be the function from Lemma 1 with the corresponding kernel $K = K_1 + a(x)K_2$.

For all $x \in B_1$, we have

$$(15) \quad -C_0 \leq L_1 \bar{u}(x) + a(x)(L_2 \bar{u}(x) - L_4 \bar{u}(x)) \leq C_0.$$

Lemma 2 implies the existence of $w(x_1)$ such that

$$(16) \quad \begin{aligned} L_1 w + a(x)(L_2 w - L_4 w) + L_3 w &\geq C_0, \quad |x_1| \leq 1, \\ L_1 w + a(x)(L_2 w - L_4 w) - L_3 w &\leq -C_0, \quad |x_1| \leq 1. \end{aligned}$$

We define $u(x) = \bar{u}(x_1) + w(x_1)$. By (15) and (16), we have

$$\begin{aligned} L_1 u + a(x)(L_2 u - L_4 u) + L_3 u &\geq 0, \\ L_1 u + a(x)(L_2 u - L_4 u) - L_3 u &\leq 0. \end{aligned}$$

if $|x_1| \leq 1$. By the intermediate value theorem, there is a $c(x) \in [-1, 1]$ so that

$$L_1 u + a(x)(L_2 u - L_4 u) - c(x)L_3 u \leq 0,$$

which implies $Lu(x) = 0$ in B_1 by (6).

Since w is also monotonically increasing in x_1 , the oscillation $\omega_{B_r} u \geq \omega_{B_r} u_r = 1$, and u breaks the modulus of continuity η . \square

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